







# Expressing general constitutive models using algorithmic automatic differentiation in DOLFINx

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#### Expressing solid mechanics problems in UFL

Weak form of an equilibrium problem in solid mechanics:

$$F(\boldsymbol{u};\boldsymbol{v}) = \int_{\Omega} \boldsymbol{\sigma}(\boldsymbol{u}) \cdot \boldsymbol{\varepsilon}(\boldsymbol{v}) d\boldsymbol{x} = 0, \quad \forall \boldsymbol{v} \in V,$$
 (1)

where  $\boldsymbol{\varepsilon}(\boldsymbol{v}) = \frac{1}{2}(\nabla \boldsymbol{v} + (\nabla \boldsymbol{v})^T).$ 

Expressing the form via UFL:

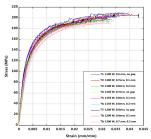
#### **UFL** limitations

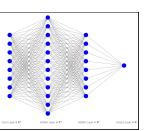
What if  $\sigma(u)$  is not expressible via analytical formulas, i.e. in UFL?

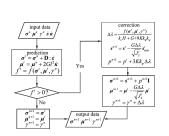
#### Issue: UFL is limited to express any constitutive model

Ways to express constitutive models  $\sigma = \sigma(u)$  in UFL:

- lacksquare  $\sigma=$  analytical formula
- $imes \sigma =$  experimental data
- $\times$   $\sigma$  = neural network
- $\times \sigma = \text{algorithm}$
- X and so on...







$$egin{aligned} rac{\partial f_{arepsilon^{
m cl}}}{\partial \Delta \, arepsilon^{
m cl}} &= oxsdel{f I} + \Delta \lambda rac{\partial n}{\partial \Delta ar{e}} \ rac{\partial f_{arepsilon^{
m cl}}}{\partial \Delta \, p} &= oldsymbol{n} \ rac{\partial f_p}{\partial \Delta \, ar{e}^{
m cl}} &= E^{-1} oldsymbol{n}_F : oldsymbol{f C} \ rac{\partial f_p}{\partial \Delta \, p} &= 0 \end{aligned}$$

#### Solution: expressing any constitutive model

#### We propose a framework that:

- ullet extends FEniCSx/DOLFINx and allows to use **any** 3rd-party library to define constitutive models  $\sigma(u)$  as a part of weak problems;
- uses NumPy arrays to pass data between DOLFINx and external libraries;
- is based on two concepts: external operator<sup>1</sup> and automatic differentiation;

<sup>&</sup>lt;sup>1</sup>Nacime Bouziani and David A. Ham. *Escaping the abstraction: a foreign function interface for the Unified Form Language [UFL]*. 2021. arXiv: 2111.00945 [cs.MS].

# What is an external operator?

An **external operator**<sup>2</sup>  $N(\cdot)$  is a *symbolic* UFL object that

- is defined by a computer program, not by an analytical formula,
- differentiable and it's derivative  $\partial N(\cdot)$  is another **external operator**.

Example:

$$F = F(u, N(u); v), \quad u, v \in V, \tag{2}$$

The Gâteaux derivative of F with respect to u in the direction  $\hat{u} \in V$ :

$$J = \mathfrak{D}_{u}F(\hat{u}) = \partial_{u}F(\hat{u}) + \partial_{N}F(\partial_{u}N(\hat{u})), \quad u, v, \hat{u} \in V,$$
(3)

where  $\mathfrak{D}_u\{\cdot\}(\hat{u})$  and  $\partial_u\{\cdot\}(\hat{u})$  are respectively total and partial Gâteaux derivatives with respect to operand u in the direction  $\hat{u}$ .

<sup>&</sup>lt;sup>2</sup>Nacime Bouziani and David A. Ham. *Escaping the abstraction: a foreign function interface for the Unified Form Language [UFL]*. 2021. arXiv: 2111.00945 [cs.MS].

#### External operators in solid mechanics

The constitutive model  $\sigma(u) := \sigma(\varepsilon(u))$  is an **external operator** acting on the strain tensor  $\varepsilon(u)$ .

To create an external operator in the FEniCSx environment, you need to use the class FEMExternalOperator of our framework:

```
1 sigma = FEMExternalOperator(
2    epsilon(u), # operand
3    function_space=S, # quadrature space
4    external_function=sigma_external # a Python function
5 )
```

The derivative  $\frac{d\sigma(\varepsilon(\boldsymbol{u}))}{d\varepsilon} = \boldsymbol{C}_{tang}(\varepsilon(\boldsymbol{u}))$ , the *consistent tangent stiffness matrix*, is another external operator and created via the UFL's automatic differentiation.

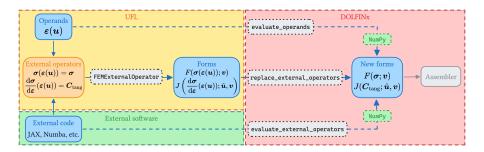
# How to define external operators

```
def sigma_impl(deps: np.ndarray) -> np.ndarray:
      return sigma_values
4
  def C_tang_impl(deps: np.ndarray) -> np.ndarray:
      return C tang values
  def sigma external(
      derivatives: Tuple[int, ...]
10
  ) -> Callable[[np.ndarray], np.ndarray]:
12
      if derivatives == (0,):
13
          return sigma_impl
14
       if derivatives == (1,):
15
          return C_tang_impl
16
```

#### Escaping FEniCSx environment

Functions above may contain **any** 3rd-party code working with NumPy arrays.

#### Framework workflow



#### Plasticity of Mohr-Coulomb

Let's consider the following weak problem for  $u \in V$ :

$$F(\mathbf{u}; \mathbf{v}) = \int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) \cdot \boldsymbol{\varepsilon}(\mathbf{v}) d\mathbf{x} - F_{\text{ext}}(\mathbf{v}) = 0, \quad \forall \mathbf{v} \in V,$$
 (4)

$$f(\boldsymbol{\sigma}) \le 0, \tag{5}$$

Non-associated Mohr-Coulomb plasticity with apex smoothing:

$$h(\boldsymbol{\sigma}, \alpha) = \frac{I_1(\boldsymbol{\sigma})}{3} \sin \alpha + \sqrt{J_2(\boldsymbol{\sigma})K^2(\alpha) + a^2(\alpha)\sin^2 \alpha} - c\cos \alpha, \tag{6}$$

$$f(\boldsymbol{\sigma}) = h(\boldsymbol{\sigma}, \phi), \text{ - yield surface}$$
 (7)

$$g(\sigma) = h(\sigma, \psi)$$
, - plastic potential (8)

where  $J_2(\boldsymbol{\sigma}) = \frac{1}{2} \boldsymbol{s} \cdot \boldsymbol{s}$  is the second invariant of the deviatoric part  $\boldsymbol{s}$  of the stress tensor and  $F_{\text{ext}}$  represents external forces.

How to solve 4-5: apply a return-mapping procedure, a numerical algorithm.

#### Reutrn-mappping procedure

Constitutive equations in plasticity:

$$\begin{cases}
\mathbf{r}_{\mathbf{g}}(\boldsymbol{\sigma}_{n+1}, \Delta \lambda) = \boldsymbol{\sigma}_{n+1} - \boldsymbol{\sigma}_{n} - \boldsymbol{C} \cdot (\Delta \varepsilon - \Delta \lambda \frac{\mathrm{d}\mathbf{g}}{\mathrm{d}\boldsymbol{\sigma}}(\boldsymbol{\sigma}_{n+1})) = \mathbf{0}, \\
\mathbf{r}_{\mathbf{f}}(\boldsymbol{\sigma}_{n+1}) = \mathbf{f}(\boldsymbol{\sigma}_{n+1}) = 0,
\end{cases} \tag{9}$$

**Return-mapping procedure** is a numerical algorithm solving the system 9 by following a predictor-corrector scheme for the stress tensor  $\sigma$ .

In the general case, e.g. the *Mohr-Coulomb* case, the return-mapping requires solving the nonlinear system 9 **numerically**.

# Plasticity problems in FEniCSx via JAX

**Our solution**: implementation of the *return-mapping* procedure using external operators via package JAX

JAX is a high-level library for automatic differentiation (AD) and numerical computing, which also supports the just-in-time compilation (JIT) feature.

# Automatic differentiation (AD) in solid mechanics via JAX

Consistent tangent stiffness matrix

$$\frac{\mathrm{d}\sigma(\varepsilon(\mathbf{u}))}{\mathrm{d}\varepsilon} = \mathbf{C}_{\mathsf{tang}}(\varepsilon(\mathbf{u})) \tag{10}$$

Implementation via JAX:

```
1 def sigma_impl(deps: np.ndarray) -> np.ndarray:
2     ...
3     "<return-mapping algorithm>"
4     ...
5     return sigma_
6
7 C_tang_impl = jax.jacfwd(sigma_impl)
```

The function  $C_{tang_impl}$  evaluate the consistent tangent stiffness matrix  $C_{tang} = \frac{d\sigma}{d\varepsilon}$  exactly at Gauss points.

# Mohr-Coulomb plasticity via JAX

Constitutive equations in plasticity:

$$\begin{cases}
\mathbf{r}_{\mathbf{g}}(\boldsymbol{\sigma}_{n+1}, \Delta \lambda) = \boldsymbol{\sigma}_{n+1} - \boldsymbol{\sigma}_{n} - \boldsymbol{C} \cdot (\Delta \varepsilon - \Delta \lambda \frac{\mathrm{d} \mathbf{g}}{\mathrm{d} \boldsymbol{\sigma}}(\boldsymbol{\sigma}_{n+1})) = \mathbf{0}, \\
\mathbf{r}_{f}(\boldsymbol{\sigma}_{n+1}) = f(\boldsymbol{\sigma}_{n+1}) = 0.
\end{cases}$$
(11)

Inner Newton loop:

```
1 def return_mapping(
2    deps_local: np.ndarray,
3    sigma_n_local: np.ndarray
4 ) -> ...:
5    ...
6    ... = jax.lax.while_loop(...)
7    ...
8    return sigma_local, (sigma_local, niter_total, yielding, norm_res,
9    dlambda)
```

The program that evaluates stress  $\sigma$  AND  $m{C}_{\mathsf{tang}} = \frac{\mathrm{d} m{\sigma}}{\mathrm{d} m{\varepsilon}}$  exactly at Gauss points:

```
1 dsigma_ddeps = jax.jacfwd(return_mapping, has_aux=True)
```

#### Mohr-Coulomb plasticity via JAX: vectorization

Defining the external operator and its derivative through the **vectorization** over quadrature points (via jax.vmap):

```
1 dsigma_ddeps_vec = jax.jit(jax.vmap(dsigma_ddeps, in_axes=(0, 0)))
```

#### Globally:

```
1 def C_tang_impl(deps: np.ndarray) -> Tuple[np.ndarray, np.ndarray]:
2    deps_ = deps.reshape((-1, 6))
3    sigma_n_ = sigma_n.x.array.reshape((-1, 6))
4
5    (C_tang_global, state) = dsigma_ddeps_vec(deps_, sigma_n_)
6    sigma_global, ... = state
7    ...
8    return C_tang_global.reshape(-1), sigma_global.reshape(-1)
```

# Plasticity before AD and JAX

Constitutive equations in plasticity:

$$\begin{cases}
\mathbf{r}_{\mathbf{g}}(\boldsymbol{\sigma}_{n+1}, \Delta \lambda) = \boldsymbol{\sigma}_{n+1} - \boldsymbol{\sigma}_{n} - \boldsymbol{C} \cdot (\Delta \varepsilon - \Delta \lambda \frac{\mathrm{d}\mathbf{g}}{\mathrm{d}\boldsymbol{\sigma}}(\boldsymbol{\sigma}_{n+1})) = \mathbf{0}, \\
\mathbf{r}_{\mathbf{f}}(\boldsymbol{\sigma}_{n+1}) = \mathbf{f}(\boldsymbol{\sigma}_{n+1}) = 0,
\end{cases} (12)$$

where  $g(\sigma) = h(\sigma, \psi)$  is a plastic potential.

Instead of

we had to manually derive<sup>3</sup>:

$$\mathbf{n} = \frac{\partial G_F}{\partial I_1} \mathbf{I} + \left( \frac{\partial G_F}{\partial J_2} + \frac{\partial G_F}{\partial \theta} \frac{\partial \theta}{\partial J_2} \right) \sigma^{\mathbf{D}} + \frac{\partial G_F}{\partial \theta} \frac{\partial \theta}{\partial J_3} J_3(\sigma^{\mathbf{D}})^{-1} : \mathbf{P}^{\mathbf{D}},$$

<sup>&</sup>lt;sup>3</sup>Thomas Helfer et al. *Invariant-based implementation of the Mohr-Coulomb elasto-plastic model in OpenGeoSys using MFront*. TFEL/MFront. URL: https://thelfer.github.io/tfel/web/MohrCoulomb.html (visited on 05/10/2024).

# Plasticity before AD and JAX

#### Instead of

```
1 drdx = jax.jacfwd(r)
2 j = drdx(x_local, deps_local, sigma_n_local)
3 ...
4 dsigma_ddeps = jax.jacfwd(sigma_return_mapping, has_aux=True)
```

we had to manually derive:4

$$\begin{cases} \frac{\partial r_{\varepsilon^{el}}}{\partial \Delta \varepsilon^{el}} = \mathbf{I} + \Delta \lambda \frac{\partial \mathbf{n}}{\partial \Delta \varepsilon^{el}} \\ \frac{\partial r_{\varepsilon^{el}}}{\partial \Delta \rho} = \mathbf{n} \\ \frac{\partial r_p}{\partial \Delta \varepsilon^{el}} = \mathbf{E}^{-1} \mathbf{n}_F : \mathbf{C} \end{cases} - > \mathbf{j} = \begin{pmatrix} \frac{\partial r_{\varepsilon_{el}}}{\partial \Delta \varepsilon_{el}} & \frac{\partial r_{\varepsilon_{el}}}{\partial \Delta \rho} \\ \frac{\partial r_p}{\partial \Delta \varepsilon^{el}} & \frac{\partial r_p}{\partial \Delta \rho} \end{pmatrix} - > \frac{\mathbf{d} \boldsymbol{\sigma}(\boldsymbol{\varepsilon}(\boldsymbol{u}))}{\mathbf{d} \boldsymbol{\varepsilon}} = \\ - > = [(\mathbf{j}^{-1})_{i,j}] \boldsymbol{C}_{\text{elas}}, \\ i, j = 1, ..., k, \\ \boldsymbol{C}_{\text{elas}} \in M_{k \times k} \end{cases}$$

(13)

<sup>&</sup>lt;sup>4</sup>Thomas Helfer et al. *Invariant-based implementation of the Mohr-Coulomb elasto-plastic model in OpenGeoSys using MFront*. TFEL/MFront. URL: https://thelfer.github.io/tfel/web/MohrCoulomb.html (visited on 05/10/2024).

#### Verification of return-mapping procedure

Yield surface of Mohr-Coulomb with apex smoothing:

$$f(\boldsymbol{\sigma}, \phi) = \frac{I_1(\boldsymbol{\sigma})}{3} \sin \phi + \sqrt{J_2(\boldsymbol{\sigma})K^2(\phi) + a^2(\phi)\sin^2 \phi} - c\cos \phi, \tag{14}$$

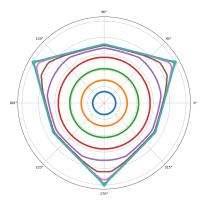


Figure: The yield surface tracing through passing several stress paths within the deviatoric plane  $(\rho, \theta)$ , where  $\rho = \sqrt{2J_2}$  and  $\theta$  is a Lode angle.

#### Verification of derivatives via Taylor test

$$R_0 = |F(\mathbf{u} + h \, \delta \mathbf{u}; \mathbf{v}) - F(\mathbf{u}; \mathbf{v})| \longrightarrow 0 \text{ at } O(h), \tag{15}$$

$$R_1 = |F(\mathbf{u} + h \, \delta \mathbf{u}; \mathbf{v}) - F(\mathbf{u}; \mathbf{v}) - J(\mathbf{u}; h \delta \mathbf{u}, \mathbf{v})| \longrightarrow 0 \text{ at } O(h^2),$$
 (16)

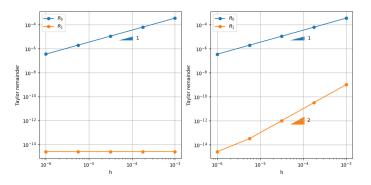


Figure: Taylor test for the form F. There are the zeroth-order  $R_0$  and the first-order  $R_1$  Taylor remainders for the elastic phase (on the left) and the plastic phase (on the right).

#### What other problems to solve?

Some examples of constitutive models where the framework can be useful:

```
• \sigma = \text{return\_mapping}(\boldsymbol{u})
```

• 
$$\sigma$$
 = another\_numerical\_algorithm( $u$ )

$$\bullet \ \sigma = FE^2(\mathbf{u})$$

• 
$$\sigma = \text{neural\_network}(u)$$

• 
$$\sigma = \text{experimental\_data}(\boldsymbol{u})$$

• 
$$\sigma = \text{call\_to\_existing\_material\_library}(\boldsymbol{u})$$

• 
$$\sigma = \text{convex\_solver}(u)$$

• 
$$\sigma = \text{surrogate\_model}(u)$$

...

#### Extending DOLFINx

The framework goes beyond the solid mechanics: It enables the support of automatic differentiation in DOFLINx via JAX!

The framework can help to integrate in DOLFINx other interesting packages and techniques:

```
■ DOLFINx + NumPy + 

■ DOLFINx + NumPy + 

■ PyTorch (neural networks), what else?, ...
```

#### Conclusion

- We implemented a framework extending DOLFINx and providing a special interface to FEniCSx users.
- This interface allows to use of any 3rd-party library to define constitutive models as a part of weak problems.
- In particular, it enables support of general automatic differentiation (AD) in FEniCSx via the **JAX** library.
- AD is a very effective and robust tool to compute derivatives defined by computer programs. This is very beneficial in the context of constitutive models.
- Just-in-time (JIT) compilation guarantees efficient implementation of a constitutive model within the framework.

Tutorials and how to use (save the link via QR code!):



=> a-latyshev.github.io/dolfinx-external-operator/

Any questions? Contact me!

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