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# Stability of functionalized platform molecules on Au(111)

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Trioxatriangulenium (TOTA) platform molecules were functionalized with methyl, ethyl, ethynyl, propynyl, and hydrogen and sublimated onto Au(111) surfaces. Low-temperature scanning tunneling microscopy data reveal that >99% of ethyl-TOTA and methyl-TOTA remain intact, whereas 60% of H-TOTA and >99% of propynyl-TOTA and ethynyl-TOTA decompose. The observed tendency toward fragmentation on Au(111) is opposite to the sequence of gas-phase stabilities of the molecules. Although Au(111) is the noblest of all metal surfaces, the binding energies of the decomposition products to Au(111) destabilize the functionalized platforms by 2 to 3.9 eV (190–370 kJ/mol) and even render some of them unstable as revealed by density functional theory calculations. Van der Waals forces are important, as they drive the adsorption of the platform molecules. *Published by AIP Publishing*. https://doi.org/10.1063/1.5059344

#### I. INTRODUCTION

The properties and functions of molecules adsorbed to metals crucially depend on the adsorption geometry and the electronic coupling to the substrate. Chemical anchor groups consisting of a molecular subunit or just a single atom are used to determine these parameters for a class of molecules. 1-15 Recently, platform molecules have been used as anchors to a substrate and to achieve decoupling from neighbor molecules. 16,17 Extended aromatic platforms are particularly interesting because they adhere to a substrate via physisorption instead of chemisorption. With these platforms lying flat on the substrate, the functional molecule juts out into vacuum. <sup>18</sup> The electronic coupling at the molecule-surface interface is mediated by an extended  $\pi$  system. Accordingly, physisorbed platforms were employed as high-conductance contacts between a metal electrode and molecular wires <sup>19,20</sup> and Zn-porphyrin molecules.<sup>21</sup> To reduce the electronic coupling, a spacer may be inserted between the functional molecule and the platform. 17,18,22-24 Photochemical cleavage of an axial phenyl acetylene from a platform molecule at the ethanol/Au(111) interface was reported.<sup>25</sup>

While Au(111) is the noblest of all metal surfaces,  $^{26}$  nm-sized Au structures can be catalytically active.  $^{27-34}$  Here we show that even on well-ordered Au(111) surfaces, some functionalized trioxatriangulenium (TOTA) platform molecules fragment despite the fact that they physisorb.  $^{35,36}$  We functionalized TOTA platforms with methyl, ethyl, ethynyl, propynyl, and hydrogen. After sublimation onto Au(111), all methyl-TOTA and >99% ethyl-TOTA on Au(111) are intact. However,  $\approx 60\%$  of the H-TOTA molecules and almost all (>99%)

## **II. EXPERIMENTAL RESULTS**

# A. Deposition of TOTA functionalized with methyl, H, and ethynyl onto Au(111)

Functionalized platforms 1–3 (Fig. 1) were sublimated onto Au(111) at mild parameters to make sure that intact molecules were deposited (see the Appendix). During sublimation, the samples were held at ambient temperature and then investigated using scanning tunneling microscope (STM) at low temperature (T = 4.6 K).

Figures 1(a) and 1(b) present typical STM topographs recorded after sublimation of molecule 1. Pristine Au(111) areas appear low (dark). A single molecule 1 is imaged as a triangular shape with a central circular protrusion [green bordered molecule in Fig. 1(b)]. The triangle (height  $\leq 0.16$  nm) and the central protrusion (0.23 nm) correspond to platform 1 and the attached CH<sub>3</sub> moiety, respectively. The platform lies flat on Au(111) with an upright CH<sub>3</sub> group as desired. The molecules arrange into hexagonal clusters, preferably on the hcp and fcc areas of the herring-bone reconstructed Au(111) substrate. These clusters are stabilized by CH···O bonds with bond angles of  $150^{\circ}$ – $170^{\circ}$  and H–O distances between 0.27 and 0.32 nm. Although we imaged thousands of molecules,

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propynyl-TOTA and ethynyl-TOTA decompose. Intriguingly, this tendency toward fragmentation on Au(111) is opposite to the sequence of gas-phase stabilities of the molecules. Density functional theory (DFT) calculations show that the different on-surface stabilities reflect the binding energies of the products to Au(111). Dispersion interaction binds the platform molecules to Au(111) and affects the positions of the intact molecules and the decomposition products but hardly affects the stability against fragmentation.

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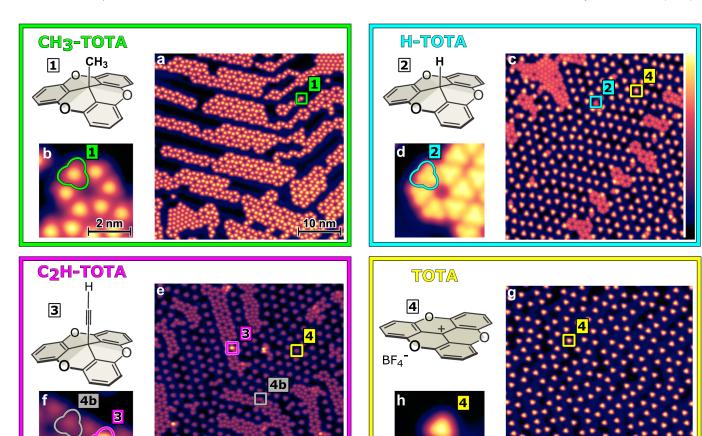


FIG. 1. [(1)-(3)] Functionalized trioxatriangulenium (TOTA) platform molecules: methyl-TOTA (1), H-TOTA (2), and ethynyl-TOTA (3) molecules. (4) Bare TOTA platform (4) with counterion BF $_4^-$ . [(a), (c), (e), and (g)] STM topographs of (1–4), respectively, on Au(111). Image width 44 nm. [(b), (d), (f), and (h)] close-up views, image width 4.4 nm, of (1–4). green, cyan, magenta, yellow, and gray symbols highlight single molecules (1–4) and (4b), respectively. The color scale shown in (c) is used throughout but corresponds to different maximum height in each image. (1–4) and (4b) exhibit maximum apparent heights of 0.23, 0.16, 0.36, 0.22, and 0.21 nm at 100 mV. (e) was recorded at -1.0 V.

we did not find a single fragmented molecule. In other words, molecule 1 remains intact on Au(111).

The deposition of molecule 2 leads to a different result [Figs. 1(c) and 1(d)]. Here two patterns with apparent heights of 0.16 and 0.21 nm [cyan and yellow squares in Figs. 1(c) and 1(d)] were found. The higher pattern (yellow) is identical to the pattern found after deposition of the bare platform 4 [Figs. 1(g) and 1(h)]. Consequently, the higher pattern is attributed to 4. In Fig. 1(c), 4 is a fragment of molecule 2. Its occurrence indicates that the bond between the central H atom, and the platform occasionally dissociates upon deposition. The lower pattern [cyan in Figs. 1(c) and 1(d)] is attributed to the intact molecule 2. In drastic contrast to the methyl-functionalized molecule 1, only  $\approx 41\%$  of more than 3000 molecules 2 were intact after deposition. The central H atom of molecule 2 is not individually resolved as expected (see, e.g., Refs. 37 and 38). Despite its similar geometry, molecule 2 appears  $\approx 0.06$  nm lower than the bare platform 1. This difference is caused by an electronic effect as discussed below.

Figures 1(e) and 1(f) present topographs recorded after deposition of molecules  $\bf 3$  (C<sub>2</sub>H–TOTA). Again, the bare

platform 4 [yellow square in Fig. 1(e)] is found as a fragment on the sample. In addition, a different pattern denoted 4b [gray markers in Figs. 1(e) and 1(f)] occurs. The pattern 4b is exclusively found after deposition of 3 with fractions varying between 5% and 80% at coverages of 0.05–0.45 molecules per 1 nm<sup>2</sup>. As discussed below, we attribute the pattern 4b to the bare platform in a different adsorption state than 4. Intact molecules 3 [magenta in Figs. 1(e) and 1(f)] appear highest (0.36 nm) because of the attached ethynyl group. However, less than 1% of 3 are intact on Au, showing that the  $C_2H$  moiety makes these molecules marginally stable on Au(111).

# B. Intermolecular interactions

The topographs in Figs. 1(c), 1(e), and 1(g) provide important additional information. In contrast to the functionalized molecules 1–3 and molecule 4b, the bare platform 4 does not cluster. The molecules 4 rather prefer to maximize their average nearest-neighbor distance d in fcc and hcp regions of the substrate. The square of the average nearest-neighbor distance,  $d^2$ , describes (up to a constant factor) the surface

area that is occupied by a single molecule. Given a long-range repulsive interaction,  $d^2$  is expected to be inversely proportional to the coverage, which determines the area that is available for a single molecule. Figure 2(a) shows that this proportionality indeed holds. Moreover, it evidences that d is not limited to specific values. This observation excludes the possibility of surface-mediated long-range interactions via Friedel oscillations of surface-state electrons, where d would change with a period of half of the Fermi wavelength, i.e., 1.8 nm for Au(111).  $^{39-41}$ 

Long-range interactions between adsorbates can also be mediated by elastic deformation of the substrate. 42-44 The herring-bone reconstruction, which is due to a stressed atomic layer, is sensitive to changes of the surface elastic energy. 45,46 However, we find that the herring-bone reconstruction is unaffected by submonolayer coverages of molecule 4. Moreover, the platform is physisorbed to the Au substrate and consequently the induction of a significant amount of strain is unlikely. Hence, an elastic surface-mediated interaction is improbable.

Next, the nearest-neighbor (NN) pair distribution functions for low coverages were investigated [Figs. 2(b) and 2(c)]. They show that the NN distances have lower and upper limits, e.g., in case of Fig. 2(b), there are no NN distances <2 or >3.5 nm. This effect along with the shape of the distributions indicates long-range Coulomb repulsion between the molecules 46 and suggests a significant charging of the adsorbed platform 4, consistent with its cationic nature in the gas phase.

Clusters of molecules **1–3** and **4b** [Figs. 1(a), 1(c), and 1(e)] exhibit the same hexagonal structure and orientation with respect to the substrate. They are stabilized by CH···O bonds between the platforms. The formation of H bonds indicates a smaller intermolecular Coulomb repulsion and consequently

less charging of the molecules, consistent with their neutral character in the gas phase. In fact, the on-surface interactions of **1–4** reflect the molecular tendency to ionize. A related switching of intermolecular interactions by attaching a ligand to a molecule has recently been reported for Fe-tetramethyl-tetraazaannulene on Au(111).<sup>47</sup>

## C. Induced dissociation and electronic states near $E_F$

As presented above, methyl-TOTA (1) is stable on Au(111) at ambient temperature, H-TOTA (2) fragments, and ethynyl-TOTA (4) is marginally stable. Moreover, the occurrence of the bare platform as fragment indicates that the bond between platform and attached moiety dissociates upon deposition. To verify that this bond is weaker than the other bonds of the adsorbed molecules, we investigated the stability of the molecules upon injecting high currents into the molecules. Indeed, injection of several nA current into the molecules 1-3 at elevated voltages ( $\gtrsim 2 \text{ V}$ ) dissociates the attached moieties from the platform but does not lead to further fragmentation of the platform. An example for such an induced dissociation is shown in Fig. 3 for molecule 1 on Au(111). Initially [Fig. 3(a)], methyl-TOTA molecules arranged in a compact island are observed. After applying a current/voltage pulse to a molecule 1 at the position indicated by a black point, a new state is obtained [Fig. 3(b)] whose characteristics match the bare platform 4. This kind of dissociation, which was repeatedly observed, is selective, i.e., no other molecules are affected. After an induced dissociation, the STM tip was usually contaminated, i.e., contacts to the bare Au substrate revealed a non-metallic tip and STM images tended to be fuzzy. We attribute this contamination to the transfer of the functional group or fragments thereof to the tip. After tip preparation by repeated dipping into the Au surface,

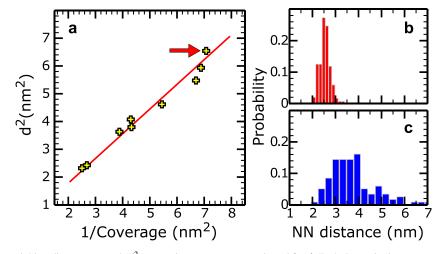


FIG. 2. (a) Average nearest-neighbor distance squared  $(d^2)$  versus inverse coverage evaluated for **4**. Each data point is an average over at least 300 molecules from large areas of a newly prepared sample. This average neglects differences in local coverage between hcp and fcc areas that are significant at low densities  $(\gtrsim 7 \text{ mm}^2 \text{ per molecule})$ . The line shows a linear fit to the data  $(d^2 = a/\gamma, \gamma)$  being the coverage in molecules per nm²). We find a = 0.89, which is close to the value expected for a hexagonal lattice  $(\sqrt{3}/2 \approx 0.87)$ , which the molecules approximately form [cf. molecules **4** in Figs. 1(c) and 1(e)]. [(b)–(c)] Nearest-neighbor (NN) pair distribution functions of **4** for coverages (b) 0.14 and (c) 0.03 molecules per nm² [Figs. 1(g) and 4(a), respectively]. The distributions were obtained from >200 molecules. Steric hindrance limits the smallest distance to  $\approx 1 \text{ nm}$ . The data point marked by an arrow in (a) corresponds to the data shown in (b).

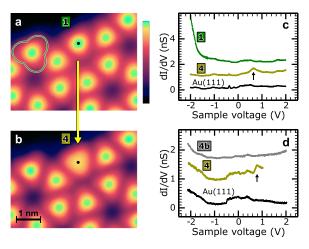


FIG. 3. [(a) and (b)] STM topographs (a) before and (b) after STM-induced dissociation of the methyl moiety from the single molecule 1 indicated by black points. [(c) and (d)] Differential conductance (dI/dV) spectra of molecule 1, 4, and 4b along with the respective background spectra of the bare Au(111) substrate. Black arrows indicate the positions of the LUMO peak of 4 when it is (c) isolated or (d) adjacent to 1. The steep rise of the conductance of 1 at  $\approx$ -1.8 eV is attributed to the HOMO. The Au(111) spectrum in (d) exhibits a rapid rise at  $\approx$ -0.5 V that is related to the edge of the surface state. The increase of the dI/dV signal in (d) for sample voltages  $\lesssim$ -1.5 eV is due to tip states. For clarity, some spectra are vertically shifted: (c) 1: +1.5 nS and 4: +0.9 nS and (d) 4b: +0.9 nS and 4: +0.3 nS. The feedback loop was opened at 0.5 nA and 2 V [1 V for 4 in (d)].

STM images could again be recorded with a clean metal tip. Figure 3(b) was actually recorded with such a re-prepared tip.

Further information about the electronic states of adsorbates can be obtained with the help of differential conductance (dI/dV) spectra, which probe the local density of states (LDOS). dI/dV spectra of 1 and 4 [Fig. 3(c)] reveal a LUMO peak at ≈0.6 eV for the bare platform 4 that is not present at molecule 1. Actually all investigated functionalized platforms do not exhibit any prominent state in the range  $|E-E_F| \lesssim 1.5$  eV. For molecule 1, a steep increase of the differential conductance is visible for energies below  $E - E_F \lesssim -1.5$  eV. This is an onset of a HOMO peak. A complete measurement of this peak at the tunneling parameters of Fig. 3(c) would lead to the aforementioned induced dissociation of 3. The dI/dV spectra of 4 in Figs. 3(c) and 3(d) reveal a further piece of information. The energy of the LUMO of 4 depends on its environment. Isolated molecules 4 show a LUMO at ≈0.6 eV [Fig. 3(c)]. An adjacent functionalized platform shifts the level away from  $E_F$ by  $\approx 0.2$  eV [Fig. 3(d)]. Polarization screening of neighboring molecules may cause this shift as observed for, e.g., perylenetetracarboxylic-dianhydride on Au(111)<sup>48</sup> and tetracene on Ag(111).<sup>49</sup>

Besides the bare platform **4**, also **4b** was identified as a fragment of **3** [Figs. 1(e) and 1(f)]. The dI/dV spectra of **4** and **4b** [Fig. 3(d)] reveal a marked difference of their electronic structures. Whereas **4** exhibits a LUMO resonance at below 1 eV, fragment **4b** has no state for  $|E-E_F| \le 2$  eV. The absence of an unoccupied state close to  $E_F$  along with the clustering of **4b** (see Sec. II B) is presumably caused by a different charge state.

#### D. Decomposition products

As to fragment 4b, several observations indicate that it is the bare platform in a different adsorption state on Au(111). 4b is found exclusively after decomposition of 3 [Figs. 1(e) and 1(f)]. Neither the sublimation of the bare platform [Figs. 1(g) and 1(h)] nor the decomposition of 2 [Figs. 1(c) and 1(d)] lead to the presence of 4b. Thus, its formation is most likely a kinetic effect of the dissociation of 3. As an intermediate, molecule 4b should be energetically less favorable than 4. Consistent with this, 4b can be converted to 4 by injecting current into the molecules at elevated voltages (≥2 V) and the reverse process could not be achieved despite various attempts at different manipulation parameters. Another intriguing observation is that the fraction of 4b after deposition of 3 drastically depends on the coverage of the molecules [Figs. 4(a) and 4(b)]. Low (high) coverages correlate with small (high) fractions of molecule **4b** [e.g., 0.05 (0.45) molecules per 1 nm<sup>2</sup> translates to  $\approx$ 5% (80%) **4b**]. When the coverage is increased from 0.29 [Fig. 1(e)] to 0.45 molecules per 1 nm<sup>2</sup> [Fig. 4(b)], the intermolecular distance of the bare platform slightly reduces by ≈0.1 nm. However, the fraction of 4b increases from 59%

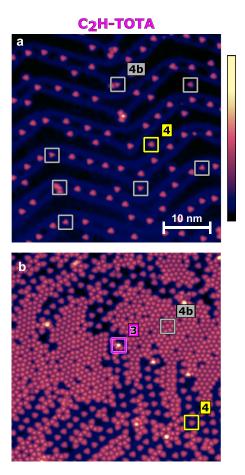


FIG. 4. STM topographs after deposition of different coverages of molecule  $\bf 3$  on Au(111): (a) 0.05 and (b) 0.45 molecules per 1 nm² (image width 44 nm). In (a), the fragments  $\bf 4b$  are at the reactive elbows of the herring-bone reconstruction and marked by a gray square, respectively. In (b), the molecules  $\bf 4b$  (gray) cluster to large islands. Molecules  $\bf 4$  (yellow) do not cluster. The color scale shown in (a) is used and corresponds to an apparent height of 3.7 nm. (b) was recorded at -1 V.

to 80%, i.e., most of the additional molecules disintegrate into molecule 4b. This coverage dependence indicates that the interaction with other molecules hinders the conversion of 4b to 4. Since 4 is significantly charged (see Sec. II B), this hindrance may be linked to the intermolecular Coulomb repulsion of 4 that renders 4, in contrast to 4b, energetically less favorable at higher coverages.

The aforementioned properties of **4b** are all similar to those of the functionalized platforms and could be explained by a moiety attached to the platform. However, there were, in contrast to the induced dissociations of the functionalized platforms (see Sec. II C), no indications of a detached moiety after conversion of **4b** to **4**. As molecule **4b** is a fragment of  $C_2H$ -TOTA, the attached moiety should comprise either one or two C atoms. Such moieties are neither chemically intuitive [high gas-phase dissociation energies of the related bonds: (i)  $HC_2$ -H:  $5.8 \, \text{eV}$  ( $560 \, \text{kJ/mol}$ ) and (ii)  $HC \equiv CH$ :  $10 \, \text{eV}$  ( $965 \, \text{kJ/mol}$ )] nor consistent with the STM topographs. Hence, fragment **4b** is most likely the bare platform in a metastable adsorption state.

While fragments **4** and **4b** are directly observed upon dissociation of **3** [Figs. 1(e) and 1(f)], the fate of the  $C_2H$  group is less clear. STM topographs reveal fuzzy areas (supplementary material, Sec. II) that may represent clusters of  $C_2H$ . The high reactivity of  $C_2H$  may also lead to the formation of gaseous molecules like  $HC_2-C_2H$  [bond formation energy -6.5 eV (627 kJ/mol) in the gas phase]<sup>50</sup> and their partial desorption at ambient temperature.

The fragmentation of molecule **2** upon adsorption on Au(111) at ambient temperature [Fig. 1(c)] may lead to H atoms on the sample—in addition to the bare platforms **4**. H atoms on Au(111) were reported to desorb at  $\approx 110$  K via the formation of  $H_2$ .<sup>51</sup> In agreement with that, we found no indications of their presence on our samples.

#### **III. CALCULATED RESULTS**

The CH<sub>3</sub> and C<sub>2</sub>H-functionalized platforms **1** and **3** are extreme cases with either exclusively intact molecules or >99% fragments on Au(111). We therefore calculated their dissociation energies with DFT assuming different reaction scenarios (see below). Moreover, we focused on the bare platform **4** as a product because we hint that **4b** may result from a cooperative effect that is difficult to tackle with existing methods. In other words, our calculations reflect the case of small coverages where most of the deposited molecules **3** fragment to **4**. To adequately describe the adsorption geometries, dispersion interactions were taken into account by using the vdW<sup>surf</sup> method (see the Appendix). <sup>52–54</sup>

# A. Dissociation energies

Reaction I represents the gas-phase dissociations of molecules 1 and 3 (Fig. 5). It turns out (Table I) that 1 and 3 are stable with dissociation energies exceeding 2 eV (190 kJ/mol), which is consistent with the observed stability toward sublimation. Moreover, in contrast to the observed instability on Au(111), the bond of the functional

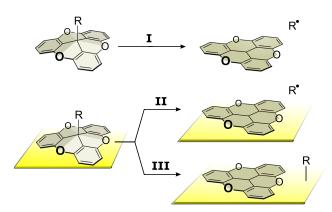


FIG. 5. Sketch of gas-phase dissociation **I** and on-surface dissociations **II** and **III** of the substituent R from the platform. In the on-surface dissociations, R either (**II**) desorbs or (**III**) adsorbs to Au(111).

group is 1 eV stronger in **3** (C<sub>2</sub>H–TOTA) than in **1**(CH<sub>3</sub>–TOTA). Reaction **I** therefore cannot explain the experimental observations.

Reaction II describes the on-surface dissociation assuming that the platform remains on Au and the functional group desorbs (Fig. 5). Interestingly, both dissociation energies (II in Table I) are  $\approx 0.45$  eV (43 kJ/mol) lower than the gas-phase values. This destabilization is mainly caused by the more favorable adsorption of the flat platform 4 ( $\approx 0.4$  eV) compared to the curved platforms with the attached functional group (Fig. 5). It is not due to the geometric deformation caused by the adsorption. Since both dissociation energies exceed 1.5 eV (144 kJ/mol) and 3 is 1 eV (96 kJ/mol) more stable than 1, II does not reflect the experimental observations.

Reaction III is an on-surface dissociation with both products remaining on Au (Fig. 5). Remarkably, III reveals that 3 (C<sub>2</sub>H–TOTA) is *not stable* on Au, while 1 (CH<sub>3</sub>–TOTA) is *stable* (Table I). The dissociation of 3 reduces the total energy by  $\approx 0.1 \text{ eV}$  (9.6 kJ/mol), whereas dissociation of 1 requires more than 0.4 eV (38 kJ/mol). In other words, the on-surface dissociation III is consistent with the experimental observation of mainly decomposed 3 and fully intact 1 on Au.

TABLE I. Dissociation energies of 1 (CH<sub>3</sub>-TOTA) and 3 (C<sub>2</sub>H -TOTA) in eV (96 kJ/mol) for calculated reaction pathways I-III. The dissociation energies were obtained as described in the section DFT of the Appendix.

	Dissociation	1	3
I	$R - TOTA \rightarrow R + TOTA$	2.43	3.77
II	$R - TOTA/Au \rightarrow R + TOTA/Au$	1.98	3.30
Ш	$R-TOTA/Au \rightarrow R/Au+TOTA/Au$	0.46	-0.10

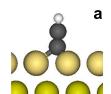
TABLE II. Adsorption energy in eV (96 kJ/mol) of the molecules 1, 3, 4 and the  $CH_3$  and  $C_2H$  radicals on Au(111). The adsorption energies were obtained as described in the section DFT of the Appendix.

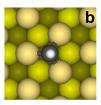
	<b>1</b> <sup>a</sup>	3 <sup>b</sup>	4 <sup>c</sup>	CH <sub>3</sub>	C <sub>2</sub> H
$E_{ad}$ (PBE + vdW)	-2.17	-2.15	-2.62	-1.5	-3.4
$E_{ad}$ (PBE)	+0.01	-0.02	-0.57	-1.3	-3.2

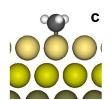
aCH<sub>3</sub>-TOTA.

 $<sup>^{</sup>b}C_{2}H$ -TOTA.

cTOTA.







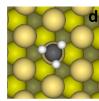


FIG. 6. Adsorption geometries of [(a) and (b)]  $C_2H$  and [(c) and (d)]  $C_3H$  and (c) show side views and (b) and (d) show top views.  $C_2H$  and  $C_3H$  are adsorbed in hollow and top sites, respectively. Yellowish: Au atoms of the first three (111) surface layers, gray: carbon, and white: hydrogen.

In addition to the dissociation pathways **I–III**, it is conceivable that the R moieties of two R–TOTA molecules directly react to R–R without the adsorption of the R moieties on the substrate. In other words, the reaction  $2 R - TOTA/Au \rightarrow R-R+2 TOTA/Au$  (denoted as **IV** below) might occur *without* reaction **III** as an intermediate step. However, the dissociation energies of reaction **IV**, 0.10 and 0.06 eV for R equal to  $C_2H_2$  and  $CH_3$ , respectively, are almost identical. As a consequence, reaction **IV** (without reaction **III** as an intermediate step) cannot explain the drastically different stabilities of molecules  $C_2H_2$ –TOTA (almost all dissociated) and  $CH_3$ –TOTA (all intact) on Au(111), in contrast to reaction **III**.

# B. Adsorption energies and geometries

The key reason for the different stabilities on Au(111) are the adsorption energies of  $C_2H$  and  $CH_3$  (Table II). The former is 1.9 eV (180 kJ/mol) more favorable than that of  $CH_3$ . Along with the lower adsorption energy of the bare platform 4, this energy difference is enough to render 3 unstable whereas 1 is not. Hence, the Au(111) surface is not inert enough for 3 to remain stable. Consistent with this interpretation, we observed that 3 dissociates on the more reactive Ag(111) surface, too (supplementary material, Sec. III).

Figure 6 shows the optimized geometries of C<sub>2</sub>H and CH<sub>3</sub> on Au(111). C<sub>2</sub>H adsorbs upright at fcc-hollow sites [Figs. 6(a) and 6(b)], whereas CH<sub>3</sub> is located on top of a surface Au atom [Figs. 6(c) and 6(d)]. The calculated Au–C distances, 0.221 nm for C<sub>2</sub>H and 0.212 nm for CH<sub>3</sub>, are typical of covalent Au–C bonds. The strong binding of C<sub>2</sub>H reflects the fact that C<sub>2</sub>H binds to three Au atoms rather than one in the case of CH<sub>3</sub>. Interestingly, CH<sub>3</sub> is marginally stable or unstable at the hollow sites, as indicated by the relaxation to the top site in the DFT calculations. We attribute this effect to the close proximity of the H atoms in CH<sub>3</sub> and the Au(111) for the expected C–Au bond distance of 0.21 nm at the hollow sites. Details of the calculated adsorption geometries of the platforms 1, 3, and 4 are presented in the supplementary material (Sec. VII).

# C. Role of dispersion interaction

The decomposition of several molecules and the relative stabilities of intermediates in on-surface reactions on Au have recently been attributed to the strong van der Waals (vdW) interaction.<sup>56–60</sup> We therefore investigated the role of dispersion for dissociation **II** and **III** by calculating the related adsorption energies with and without inclusion of

vdW interaction (Table II), namely, with the PBE + vdW functional and the pure PBE (Perdew-Burke-Ernzerhof) functional (see Appendix, Subsection 4). The adsorption energies of 1, 3, and 4 increase by ≈2.2 eV with PBE compared to PBE + vdW, reflecting strong dispersion interaction with Au(111) and its crucial role for the adsorption process. The adsorption energies of 1 and 3 may contribute to overcoming the energy barrier of dissociation process III. However, the calculated adsorption energies of 1 and 3 are almost identical and therefore cannot explain the observed stability difference. Similarly, dispersion interaction has little influence on the dissociation energies when the correct adsorption geometries are used. The energy differences relevant for dissociation paths II and III, namely, between 4 on one hand and 1 and 3 on the other ( $\approx 0.4 \text{ eV}$ ) and between CH<sub>3</sub> and C<sub>2</sub>H ( $\approx 1.9 \text{ eV}$ ), are almost unaffected when changing from PBE + vdW to PBE. Hence, the dispersion interaction plays an indirect role in the reactions. It defines the position of the adsorbates on the Au(111) surface but cancels out from the reaction energies. We emphasize that vdW interactions are crucial as not taking them into account, i.e., only using the pure PBE functional, leads to incorrect adsorption geometries and unreasonable dissociation energies.

# IV. DISCUSSION

Experimentally we observed that molecule 3 fragments with a very high probability. If a reaction barrier is present, it obviously can be overcome at ambient temperature. In any event, according to our DFT results, dissociation lowers the energy of the system. As to molecule 1, our calculations show that dissociation is energetically unfavorable in agreement with its observed stability. An activation barrier toward dissociation may add to this stability but cannot be evaluated from the available data.

The pathway to dissociation is a multi-step process, and it is not realistic to model the complete dissociation with existing *ab initio* methods. <sup>61</sup> Nonetheless, it is interesting to speculate about the mechanisms of fragmentation and molecular assembly. Dissociation **III** may require that the functional group is oriented toward the surface. Disregarding a hypothetical steering effect of the surface during the impact of the molecules, approximately half of the molecules are expected to impinge with the platform oriented toward the surface. The observations for **3** [Fig. 1(e)] imply that this orientation during impact is no obstacle to fragmentation. However, the barrier for flipping **3** on a flat area is expected to exceed 1 eV (96 kJ/mol) because of its large adsorption energy (Table II). Diffusion of

the molecules to step edges may facilitate a reorientation. It is also conceivable that diffusing Au atoms may attack the ligand bond of 3.

Our interpretation of the fragmentation data is consistent with the results from other functionalized TOTA platforms. We synthesized propynyl-TOTA, which may be viewed as methyl-TOTA with an inserted ethynyl spacer. As expected, we find propynyl-TOTA to be only marginally stable on Au(111) with <1% molecules remaining intact (see supplementary material, Sec. IV). The stability of ethyl-TOTA is also in line with our model. The gas-phase dissociation energy of ethyl-TOTA and the binding energy of ethyl to Au(111) are expected to be similar to those of methyl. Indeed, in our experiments, we found almost only (>99%) intact ethyl-TOTA molecules on Au(111) (see supplementary material, Sec. IV).

# V. CONCLUSION

The interaction of closed-shell platform molecules with the inert Au(111) surface surprisingly renders a rather stable  $sp_3$  C bond within a molecule unstable. The stability of the molecules depends drastically on their functional moiety. Modeling shows that the binding energies of the fragments to Au(111) destabilize the  $sp_3$  C bond between the functional moieties and the platform by 2-3.9 eV (190-370 kJ/mol) and essentially determine the molecular stability. Dispersion interaction plays an important indirect role, as it binds the platforms to the Au(111) surface, but it cancels out from the reaction energies. Moreover, the results show that the thermal energy at ambient temperature and the impact of the molecules are sufficient to overcome kinetic limitations.

#### SUPPLEMENTARY MATERIAL

See Supplementary material for details of synthesis and calculations; calculated MODOS of adsorbed molecules 1 and 4; topographs of propynyl-TOTA, ethyl-TOTA, and ethynyl fragments on Au(111); topographs of  $\mathbf{1}$  and  $\mathbf{4}$  on Ag(111); calculated adsorption geometries of 1, 3, and 4.

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# APPENDIX: EXPERIMENTAL AND THEORETICAL **DETAILS**

#### 1. Synthesis of the molecules

Functional groups were attached to the cationic TOTA platform 4 by polar C-C and C-H bond formation as described in the supplementary material in Sec. I. The powders 1-3 were purified by sublimation at  $\approx 0.5$  Pa and  $\approx 100$  °C.

# 2. Sample preparation

Submonolayers of TOTA<sup>+</sup>BF<sub>4</sub> salt and the molecules 1-3 were sublimated onto flat and clean Au(111) held at

ambient temperature. The sublimation parameters for molecules 1-3 were almost identical ( $<10^{-7}$  Pa and  $\approx 70$  °C) and milder than during the aforementioned purification by sublimation to make sure that intact molecules were deposited. The time between sublimation (<2 min) and cool down of the sample to ≈10 K for STM measurments did not exceed 25 min.

#### 3. STM

STM images were recorded at 4.6 K, <10<sup>-9</sup> Pa, low constant currents (20-40 pA), and low positive sample voltages (100–130 mV). All STM topographs were acquired using the constant-current mode.

## 4. DFT

Electronic structure calculations were performed by using DFT with non-empirical exchange-correlation functional of Perdew-Burke-Ernzerhof<sup>62</sup> (PBE) and the vdW<sup>surf</sup> method to model screened van der Waals interactions for atoms and molecules on surfaces<sup>52-54</sup> as implemented in the FHI-aims code<sup>63</sup> denoted PBE + vdW above.

The dissociation energies of reaction I (Table I) were obtained in two steps:  $H + R - TOTA \rightarrow H - R + TOTA$  $\rightarrow$  H + R + TOTA. The first step was calculated with DFT + vdW. For the second step, experimental values of the dissociation energies of methane and ethyne measured at 298 K, 4.55 eV, and 5.78 eV, were used.<sup>64</sup> The dissociations II (Table I) were calculated in the same way but with R – TOTA and TOTA adsorbed on Au(111). All data for the reaction III (Table I) were obtained with DFT + vdW calculations. Finally, the dissociation energies of reaction IV (see Sec. III A) were determined in two steps:  $2 R - TOTA/Au \rightarrow 2$  $(R + TOTA/Au) \rightarrow R - R + 2 TOTA/Au$ . The first step is reaction II for two isolated R - TOTA molecules on Au and was calculated as mentioned above. For the second step, experimental values of the bond formation energies of H<sub>3</sub>C-CH<sub>3</sub> and H<sub>2</sub>C-CH<sub>2</sub> measured at ambient temperature, -3.91 eV<sup>64</sup> and -6.51 eV,<sup>50</sup> were used.

The adsorption energies of 1, 3, and 4 (Table II) were directly calculated with the PBE + vdW and PBE functionals. The adsorption energies of the radicals were extracted from Table I as difference of the dissociation energies of the reactions II and III.

Further calculation details are given in the supplementary material in Sec. V.

<sup>&</sup>lt;sup>1</sup>K. Ichimura, S. K. Oh, and M. Nakagawa, Science 288, 1624 (2000).

<sup>&</sup>lt;sup>2</sup>E. Galoppini, W. Guo, W. Zhang, P. G. Hoertz, P. Qu, and G. J. Meyer, J. Am. Chem. Soc. 124, 7801 (2002).

<sup>&</sup>lt;sup>3</sup>H. Jian and J. M. Tour, J. Org. Chem. **68**, 5091 (2003).

<sup>&</sup>lt;sup>4</sup>Q. Li, A. V. Rukavishnikov, P. A. Petukhov, T. O. Zaikova, C. Jin, and J. F. W. Keana, J. Org. Chem. 68, 4862 (2003).

<sup>&</sup>lt;sup>5</sup>T. Kitagawa, Y. Idomoto, H. Matsubara, D. Hobara, T. Kakiuchi, T. Okazaki, and K. Komatsu, J. Org. Chem. 71, 1362 (2006).

<sup>&</sup>lt;sup>6</sup>S. Wagner, F. Leyssner, C. Kördel, S. Zarwell, R. Schmidt, M. Weinelt, K. Rück-Braun, M. Wolfad, and P. Tegeder, Phys. Chem. Chem. Phys. 11, 6242 (2009).

<sup>&</sup>lt;sup>7</sup>U. Siemeling, C. Bruhn, F. Bretthauer, M. Borg, F. Träger, F. Vogel, W. Azzam, M. Badin, T. Strunskus, and C. Wöll, Dalton Trans. 2009, 8593.

<sup>&</sup>lt;sup>8</sup>K. Moth-Poulsen and T. Bjørnholm, Nat. Nanotechnol. 4, 551 (2009).

<sup>&</sup>lt;sup>9</sup>F. von Wrochem, D. Gao, F. Scholz, H.-G. Nothofer, G. Nelles, and J. M. Wessels, Nat. Nanotechnol. 5, 618 (2010).

- <sup>10</sup>A. Schramm, C. Stroh, K. Dössel, M. Lukas, M. Fischer, F. Schramm, O. Fuhr, H. v. Löhneysen, and M. Mayor, Eur. J. Inorg. Chem. 2013, 70.
- <sup>11</sup>W. Li, L. G. C. Rego, F.-Q. Bai, J. Wang, R. Jia, L.-M. Xie, and H.-X. Zhang, J. Phys. Chem. Lett. 5, 3992 (2014).
- <sup>12</sup>E. Leary, A. L. Rosa, M. T. González, G. Rubio-Bollinger, N. Agraït, and N. Martín, Chem. Soc. Rev. 44, 920 (2015).
- <sup>13</sup>R. Frisenda, S. Tarkuç, E. Galán, M. L. Perrin, R. Eelkema, F. C. Grozema, and H. S. J. van der Zant, Beilstein J. Nanotechnol. 6, 1558 (2015).
- <sup>14</sup>Y. Ie, K. Tanaka, A. Tashiro, S. K. Lee, H. R. Testai, R. Yamada, H. Tada, and Y. Aso, J. Phys. Chem. Lett. 6, 3754 (2015).
- <sup>15</sup>C. Schuschke, M. Schwarz, C. Hohner, T. N. Silva, L. Fromm, T. Döpper, A. Görling, and J. Libuda, J. Phys. Chem. Lett. 9, 1937 (2018).
- <sup>16</sup>P. Du, M. Jaouen, A. Bocheux, C. Bourgogne, Z. Han, V. Bouchiat, D. Kreher, F. Mathevet, C. Fiorini-Debuisschert, F. Charra, and A.-J. Attias, Angew. Chem., Int. Ed. 53, 10060 (2014).
- <sup>17</sup>M. Valášek and M. Mayor, Chem.–Eur. J. **23**, 13538 (2017).
- <sup>18</sup>B. Baisch, D. Raffa, U. Jung, O. M. Magnussen, C. Nicolas, J. Lacour, J. Kubitschke, and R. Herges, J. Am. Chem. Soc. **131**, 442 (2009).
- <sup>19</sup>Z. Wei, X. Wang, A. Borges, M. Santella, T. Li, J. K. Sørensen, M. Vanin, W. Hu, Y. Liu, J. Ulstrup, G. C. Solomon, Q. Chi, T. Bjørnholm, K. Nørgaard, and B. W. Laursen, Langmuir 30, 14868 (2014).
- <sup>20</sup>T. Jasper-Tönnies, A. Garcia-Lekue, T. Frederiksen, S. Ulrich, R. Herges, and R. Berndt, Phys. Rev. Lett. **119**, 066801 (2017).
- <sup>21</sup>N. Hauptmann, L. Groß, K. Buchmann, K. Scheil, C. Schütt, F. L. Otte, R. Herges, C. Herrmann, and R. Berndt, New J. Phys. 17, 013012 (2015).
- <sup>22</sup>F. Matino, G. Schull, F. Köhler, S. Gabutti, M. Mayor, and R. Berndt, Proc. Natl. Acad. Sci. U. S. A. 108, 961 (2010).
- <sup>23</sup> N. L. Schneider, F. Matino, G. Schull, S. Gabutti, M. Mayor, and R. Berndt, Phys. Rev. B **84**, 153403 (2011).
- <sup>24</sup>S.-E. Zhu, Y.-M. Kuang, F. Geng, J.-Z. Zhu, C.-Z. Wang, Y.-J. Yu, Y. Luo, Y. Xiao, K.-Q. Liu, Q.-S. Meng, L. Zhang, S. Jiang, Y. Zhang, G.-W. Wang, Z.-C. Dong, and J. G. Hou, J. Am. Chem. Soc. 135, 15794 (2013).
- <sup>25</sup>S. Yokoyama, T. Hirose, and K. Matsuda, Chem. Lett. **44**, 1616 (2015).
- <sup>26</sup>B. Hammer and J. K. Norskov, Nature **376**, 238 (1995).
- <sup>27</sup> A. Sanchez, S. Abbet, U. Heiz, W.-D. Schneider, H. Häkkinen, R. N. Barnett, and U. Landman, J. Phys. Chem. A 103, 9573 (1999).
- <sup>28</sup>M. Haruta, Chem. Rec. **3**, 75 (2003).
- <sup>29</sup>M. Chen and D. Goodman, Science **306**, 252 (2004).
- <sup>30</sup> M. Chen, Y. Cai, Z. Yan and D. W. Goodman, J. Am. Chem. Soc. **128**, 6341 (2006).
- <sup>31</sup>A. S. K. Hashmi and G. J. Hutchings, Angew. Chem., Int. Ed. 45, 7896 (2006).
- <sup>32</sup>B. Hvolbæk, T. V. Janssens, B. S. Clausen, H. Falsig, C. H. Christensen, and J. K. Nørskov, Nano Today 2, 14 (2007).
- <sup>33</sup>H.-T. Chen, J.-G. Chang, S.-P. Ju, and H.-L. Chen, J. Phys. Chem. Lett. 1, 739 (2010).
- <sup>34</sup>J. Gong, Chem. Rev. **112**, 2987 (2011).
- <sup>35</sup>For small molecules the binding energy was often used to define physisorption and chemisorption. This approach is not useful for largish molecule where dispersion interaction may be significant. For example, the calculated adsorption energy of 4-mercaptobenzoic acid on Au(111) of 1.75 eV is mainly due to dispersion interaction (1.2 eV).<sup>36</sup>
- <sup>36</sup>N. Hauptmann, R. Robles, P. Abufager, N. Lorente, and R. Berndt, J. Phys. Chem. Lett. 7, 1984 (2016).

- <sup>37</sup>P. Liljeroth, J. Repp., and G. Meyer, Science **317**, 1203 (2007).
- <sup>38</sup>S. Pan, Q. Fu, T. Huang, A. Zhao, B. Wang, Y. Luo, J. Yang, and J. Hou, Proc. Natl. Acad. Sci. U. S. A. **106**, 15259 (2009).
- <sup>39</sup> J. Repp, F. Moresco, G. Meyer, K.-H. Rieder, P. Hyldgaard, and M. Persson, Phys. Rev. Lett. 85, 2981 (2000).
- <sup>40</sup>S. Lukas, G. Witte, and C. Wöll, Phys. Rev. Lett. **88**, 028301 (2001).
- <sup>41</sup>N. Knorr, H. Brune, M. Epple, A. Hirstein, M. A. Schneider, and K. Kern, Phys. Rev. B 65, 115420 (2002).
- <sup>42</sup>K. Lau and W. Kohn, Surf. Sci. **75**, 69 (1978).
- <sup>43</sup>H. Brune, M. Giovannini, K. Bromann, and K. Kern, Nature **394**, 451 (1998).
- <sup>44</sup>K. Pohl, M. C. Bartelt, J. de la Figuera, N. C. Bartelt, J. Hrbek, and R. Q. Hwang, Nature **397**, 238 (1999).
- <sup>45</sup>S. Narasimhan and D. Vanderbilt, Phys. Rev. Lett. **69**, 1564 (1992).
- <sup>46</sup>I. Fernandez-Torrente, S. Monturet, K. J. Franke, J. Fraxedas, N. Lorente, and J. I. Pascual, Phys. Rev. Lett. 99, 176103 (2007).
- <sup>47</sup>T. Knaak, T. G. Gopakumar, B. Schwager, F. Tuczek, R. Robles, N. Lorente, and R. Berndt, J. Am. Chem. Soc. 138, 7544 (2016).
- <sup>48</sup>J. Kröger, H. Jensen, R. Berndt, R. Rurali, and N. Lorente, Chem. Phys. Lett. 438, 249 (2007).
- <sup>49</sup>S. Soubatch, C. Weiss, R. Temirov, and F. S. Tautz, Phys. Rev. Lett. **102**, 177405 (2009).
- <sup>50</sup>S. W. Benson, J. Chem. Educ. **42**, 502 (1965).
- <sup>51</sup>M. Pan, D. W. Flaherty, and C. B. Mullins, J. Phys. Chem. Lett. 2, 1363 (2011).
- <sup>52</sup> V. G. Ruiz, W. Liu, E. Zojer, M. Scheffler, and A. Tkatchenko, Phys. Rev. Lett. **108**, 146103 (2012).
- <sup>53</sup>R. J. Maurer, V. G. Ruiz, and A. Tkatchenko, J. Chem. Phys. **143**, 102808 (2015).
- <sup>54</sup>V. G. Ruiz, W. Liu, and A. Tkatchenko, Phys. Rev. B **93**, 035118 (2016).
- <sup>55</sup>G. Bistoni, P. Belanzoni, L. Belpassi, and F. Tarantelli, J. Phys. Chem. A 120, 5239 (2016).
- <sup>56</sup>J. C. F. Rodriguez-Reyes, C. G. F. Siler, W. Liu, A. Tkatchenko, C. M. Friend, and R. J. Madix, J. Am. Chem. Soc. **136**, 13333 (2014).
- <sup>57</sup>S. Karakalos, Y. Xu, F. C. Kabeer, W. Chen, J. C. F. Rodríguez-Reyes, A. Tkatchenko, E. Kaxiras, R. J. Madix, and C. M. Friend, J. Am. Chem. Soc. 138, 15243 (2016).
- <sup>58</sup>S. Ossinger, H. Naggert, L. Kipgen, T. Jasper-Toennies, A. Rai, J. Rudnik, F. Nickel, L. M. Arruda, M. Bernien, W. Kuch, R. Berndt, and F. Tuczek, J. Phys. Chem. C **121**, 1210 (2017).
- <sup>59</sup>Y. Xu, W. Chen, E. Kaxiras, C. M. Friend, and R. J. Madix, J. Phys. Chem. B 122, 555 (2017).
- <sup>60</sup>F. C. Kabeer, W. Chen, R. J. Madix, C. M. Friend, and A. Tkatchenko, J. Phys. Chem. C **121**, 27905 (2017).
- <sup>61</sup>For example, the unit cell required for reliable adsorption geometries and energies contains 305 atoms (See supplementary material, Sec. V). This renders *ab initio* calculations of reaction barriers with any reliable level of accuracy extremely challenging, if not impossible.
- <sup>62</sup>J. P. Perdew, K. Burke, and M. Ernzerhof, Phys. Rev. Lett. **77**, 3865 (1996).
- <sup>63</sup>V. Blum, R. Gehrke, F. Hanke, P. Havu, V. Havu, X. Ren, K. Reuter, and M. Scheffler, Comput. Phys. Commun. 180, 2175 (2009).
- <sup>64</sup>S. J. Blanksby and G. B. Ellison, Acc. Chem. Res. **36**, 255 (2003).